

Relation between the High Density Phase and the Very-High Density Phase of Amorphous Solid Water

Nicolas Giovambattista¹, H. Eugene Stanley¹ and Francesco Sciortino²

¹*Center for Polymer Studies and Department of Physics,
Boston University, Boston, MA 02215 USA*

²*Dipartimento di Fisica and INFN Udr and Center for Statistical Mechanics and Complexity,
Universita' di Roma "La Sapienza"*

Piazzale Aldo Moro 2, I-00185, Roma, Italy

(Dated: 11 March 2004)

It has been suggested that high-density amorphous (HDA) ice is a structurally arrested form of high-density liquid (HDL) water, while low-density amorphous (LDA) ice is a structurally arrested form of low-density liquid (LDL) water. Recent experiments and simulations have been interpreted to support the possibility of a second “distinct” high-density structural state, named very high-density amorphous (VHDA) ice, questioning the LDL-HDL hypothesis. We test this interpretation using extensive computer simulations, and find that VHDA is a more stable form of HDA and that in fact VHDA should be considered as the amorphous ice of the quenched HDL.

PACS numbers:

The most common form of water in the universe is not a liquid but is a disordered solid named glassy water [1]. Depending on the glass formation route—vapor deposition, hyperquenching of the liquid, pressure induced crystal amorphization—water forms amorphous solids of quite different structure, with the density varying by as much as 40 percent between the lowest density form and the highest [2, 3]. The structural and thermodynamical properties of the different amorphous forms of glassy water have been the focus of many recent experimental [4, 5, 6, 7, 8, 9, 10, 11], numerical and theoretical studies [12, 13, 14]. However, a clear picture of the phase diagram of glassy water is still missing. This is in part because the properties of glasses change drastically with aging [8, 11] and because experiments at high pressure are difficult to perform. Glassy water can be found in at least two different forms, low-density amorphous (LDA) and high-density amorphous (HDA) ice [15, 16]. If a glass is formed by extremely rapid cooling of a liquid, then one can naturally associate the glass with the liquid state. LDA can be obtained experimentally by very fast quenching of low-density liquid (LDL) water at normal pressure [17]. However, the generation of an HDA glass via high pressure cooling of the liquid have proved an elusive goal; in fact, HDA is experimentally formed not by quenching but rather by compression either of LDA or of crystalline ice [16]. Therefore, the actual relationship of glassy water to the liquid at high pressure remains unknown. On the other side, computer simulations [18, 19] indicate that HDA is a structurally arrested form of high-density liquid (HDL) water, and therefore, HDA can be obtained by quenching of HDL at high pressure. A second “distinct” high-density structural state, named very high-density amorphous solid (VHDA) has been recently discovered. The density of VHDA is 7-8% higher than the density of HDA. VHDA is generated by heating HDA under pressure. The resulting glass does not convert back to HDA when recovered at ambient pressure at $T = 77\text{K}$.

The possibility of a VHDA phase in addition to an HDA phase raises many interesting questions [20], such as the relation between VHDA and HDA; this relation is important to elucidate the hypothesis that below a critical temperature there are two distinct phases of liquid water, LDL and HDL [18]. Which phase, HDA or VHDA, more resembles liquid water at high pressure is particularly interesting, and it appears that the structural properties of VHDA are closer to liquid water at high pressure than those of HDA [5, 12]. On the other side, the presence of two different high-density amorphous ices could imply the existence of more than one form of HDL as recently suggested by computer simulations [21].

Computer simulations offer a potentially useful tool to probe the relation between HDA, VHDA and liquid water, since the time scale of most simulations is sufficiently short that liquid and glassy states can be studied for a wider range of state points than is accessible experimentally. Here we report a set of extensive simulations that suggest that VHDA not HDA, may be considered as a physical manifestation of the quenched high pressure liquid. Further, our simulations suggest that VHDA is not a new thermodynamically distinct structure but rather VHDA results from partial annealing of the HDA structures made possible by the higher annealing temperature. Hence HDA is not stable but rather is highly metastable, relaxing to VHDA in a fashion analogous to the way that, on slow heating, glasses generated with hyperquenched methods relax to glasses generated with standard cooling rates.

Simulations offer the unique possibility of comparing (i) the glass resulting from conversion of HDA to VHDA with (ii) the glass generated by isobaric cooling of the liquid. However, a note of caution is in order, since in comparing experiments and simulations for glasses, one must carefully account for the significantly different time scales probed. State-of-the-art simulations probe time scales of 10–100 ns, so when the characteristic relaxation

time becomes comparable to this time, the system glassifies. Since the experimental homogeneous nucleation time is longer than 100 ns, glass configurations can be generated in simulations by cooling the liquid both at low and at high pressure. We shall exploit this fact in the present study.

We simulate a system of 216 molecules, using the simple point charge extended (SPC/E) model of water [22]. This model has been studied extensively, and the ρ and T dependence of structural and dynamic properties in equilibrium are known. In particular, the SPC/E model reproduces the thermodynamic anomalies of water—e.g., it produces a maximum in ρ [23]. At $\rho = 0.94 \text{ g/cm}^3$ at low T , the model describes well the LDA structure. We integrate the equation of motion using a time step of 1 fs and implement the reaction field method to account for long range forces. Results are averaged over 16 independent realizations. LDA ice configurations are generated by cooling to 77 K equilibrium liquid configurations at $\rho = 0.90 \text{ g/cm}^3$. To generate HDA configurations, we compress LDA at $T = 77 \text{ K}$ to $P > 1 \text{ GPa}$ at a rate of $5 \times 10^{-5} \text{ g/cm}^3/\text{ps}$ [14]. The HDA to VHDA transition is studied heating HDA at constant pressure from $T = 77 \text{ K}$ using a heating rate of 30 K/ns. We have used the same heating/cooling and compression/decompression rates in all our calculations.

Figure 1 compares, for a high pressure isobar, available experimental data (Fig. 1A) and simulation results (Fig. 1B). In Fig. 1A, the density ρ vs. temperature T for equilibrium liquid water is complemented with the ρ vs. T data of Ref. [4] obtained by heating HDA from 77 K up to 165 K, and then cooling back to 77 K. VHDA is the glass that results from densification of HDA during heating under high pressure. VHDA can be cycled, at constant pressure, between 77 K and 165 K without significant further density changes, leading us to hypothesize that VHDA is the more stable form of HDA. Our hypothesis is consistent with Mishima’s observation that HDA samples annealed to 130-150 K at 1-1.5 GPa are characterized by identical scattering patterns [24]. Relaxation of HDA to VHDA, and the irreversibility of such transformation upon further temperature-cycling, are facts reminiscent of slow heating of hyperquenched glasses, so the conversion of HDA to VHDA can be interpreted as a temperature-driven partial equilibration of the sample.

In simulations HDA is generated—in analogy with the experimental procedure—by compressing LDA or ice I_h at 77 K [25], while LDA is generated by cooling of the liquid at ambient pressure. Figure 1B shows ρ as HDA is heated. At about 100 K, ρ begins to increase, reaching at around 140 K a value 0.04 g/cm^3 larger than the density of HDA at 77 K. Decreasing T back to 77 K does not regenerate the HDA density, but instead ρ increases, in agreement with experimental results. The resulting denser material—which we identify with VHDA—can be cycled back and forth without significant changes in ρ .

Next we test the hypothesis that VHDA is the glass that would be generated by cooling the high pressure liquid. To this end, Fig. 2 compares ρ for isobaric cooling of the liquid with ρ for VHDA, two systems generated by *completely different thermal and pressure histories*. The ρ data display remarkable similarity, supporting the possibility that VHDA is indeed the physical realization of the glass generated by quenching the high pressure liquid. To further support this interpretation, Fig. 3 shows that the radial distribution functions of VHDA and of the glass obtained by isobaric cooling of the liquid are indeed indistinguishable.

The results presented in Figs. 2 and 3 indicate that compression of LDA at 77 K generates a system that is kinetically trapped due to the low temperature. This system relaxes to the more stable VHDA, whose structure is identical to the glass generated by quenching the high pressure supercooled liquid. In this respect, not HDA and LDA but rather VHDA and LDA should be thought of as the two distinct glassy states arising from the two distinct liquids associated with the hypothesized line of liquid-liquid phase transitions.

If indeed HDA is a partially-equilibrated glass generated by the compression technique, then one expects that VHDA should not convert to HDA under any transformation which does not involve a different “intermediate” phase (such as LDA, or a crystal phase). To confirm this expectation, we recover VHDA at $T = 77 \text{ K}$ and ambient pressure with $\rho = 1.26 \text{ g/cm}^3$. Next we isochorically heat this system to 155 K (Fig. 4). We compare the structure of the resulting system with the structure of the VHDA glass recovered at the same density at $T = 155 \text{ K}$. We find that the two glasses, which are generated from two completely different histories, are identical both thermodynamically (Fig. 4A) and structurally (Fig. 4B), and further we find that no VHDA to HDA transition takes place along the loop. Our conclusion motivates the need for a definitive experimental test that no VHDA \rightarrow HDA transition occurs. When VHDA is heated isochorically from 77 to 140 K, from a starting pressure of 0.02 GPa, the width of the x-ray pattern remains narrow like VHDA [9] (and hence does not appear to revert to HDA); however the position of its first maximum shifts (about 2 degrees) in the direction of HDA (cf. Figs. 1A’, 1B’, and 1D’ of Ref.[9]).

In summary, our simulations suggest that VHDA is not a new form of glassy water but it is the result of annealing HDA upon heating at high pressure. We also find that VHDA is the glass obtained by fast quenching of high pressure water, i.e. HDL. How HDA and VHDA behave upon decompression to normal pressure and the phase diagram of amorphous water are relevant topics that will be addressed in a separate work.

We thank NSF, MIUR Cofin 2002 and Firb and INFN Pra GenFdt for support and the Boston University Computation Center for a generous allocation of CPU time.

- [1] P. Jenniskens, S. F. Barnhak, D. F. Blake, and M. R. S. McCoustra, *J. Chem. Phys.* **107**, 1232–1241 (1997).
- [2] C. A. Angell, *Ann. Rev. Chem.* **xx**, xxx–xxx (review article in press).
- [3] P. G. Debenedetti, *J. Phys.: Condens. Matter* **15**, R1669–R1726 (2003).
- [4] T. Loerting, C. Salzmann, I. Kohl, E. Mayer, and A. Hallbrucker, *Phys. Chem. Chem. Phys.* **3**, 5355 (2001).
- [5] J. L. Finney, D. T. Bowron, A. K. Soper, T. Loerting, E. Mayer, and A. Hallbrucker, *Phys. Rev. Lett.* **89**, 205503–205506 (2002).
- [6] S. Klotz, G. Hamel, J. S. Loveday, R. J. Nelmes, M. Guthrie, and A. K. Soper, *Phys. Rev. Lett.* **89**, 285502–285505 (2002).
- [7] M. Guthrie, J. Urquidi, C. A. Tulk, C. J. Benmore, D. D. Klug, and J. Neufeind, *Phys. Rev. B* **68**, 184110–184114 (2003).
- [8] C. A. Tulk, C. J. Benmore, J. Urquidi, D. D. Klug, J. Neufeind, B. Tomberli, and P. A. Egelstaff, *Science* **297**, 1320–1323 (2002).
- [9] T. Loerting, I. Kohl, C. Salzmann, E. Mayer, and A. Hallbrucker, *J. Chem. Phys.* **116**, 3171–3174 (2002).
- [10] J. L. Finney, A. Hallbrucker, I. Kohl, A. K. Soper, and D. T. Bowron, *Phys. Rev. Lett.* **88**, 225503–225506 (2002).
- [11] O. Mishima and Y. Suzuki, *Nature* **491**, 599–603 (2002).
- [12] B. Guillot and Y. Guissani, *J. Chem. Phys.* **119**, 11740–11752 (2003).
- [13] V. V. Brazhkin and A. G. Lyapin, *J. Phys.: Condens. Matter* **15**, 6059–6084 (2003).
- [14] N. Giovambattista, H. E. Stanley, and F. Sciortino, *Phys. Rev. Lett.* **91**, 115504–115507 (2003).
- [15] O. Mishima, K. Takemura, and K. Aoki, *Science* **254**, 406 (1991).
- [16] O. Mishima, L. D. Calvert, and E. Whalley, *Nature* **314**, 76–78 (1985).
- [17] P. Brüggeller and E. Mayer, *Nature* **288**, 569–571 (1980).
- [18] P. H. Poole, F. Sciortino, U. Essmann, and H. E. Stanley, *Nature* **360**, 324–328 (1992).
- [19] P. H. Poole, U. Essmann, F. Sciortino, and H. E. Stanley, *Phys. Rev. E* **48**, 4605–4610 (1993).
- [20] D. D. Klug, *Nature* **420**, 749–750 (2002).
- [21] I. Brovchenko, A. Geiger, and A. Oleinikova, *J. Chem. Phys.* **118**, 9473 (2003).
- [22] H. J. C. Berendsen, J. R. Grigera, and T. P. Stroatsma, *J. Phys. Chem.* **91**, 6269–6271 (1987).
- [23] S. Harrington, P. H. Poole, F. Sciortino, and H. E. Stanley, *J. Chem. Phys.* **107**, 7443–7450 (1997).
- [24] O. Mishima, *Nature* **384**, 546–549 (1996).
- [25] J. S. Tse and M. L. Klein, *Phys. Rev. Lett.* **58**, 1672–1675 (1987).
- [26] L. Haar, J. S. Gallagher, and G. Kell, *NBS/NRC Steam Tables* (Hemisphere, Washington DC, 1985).

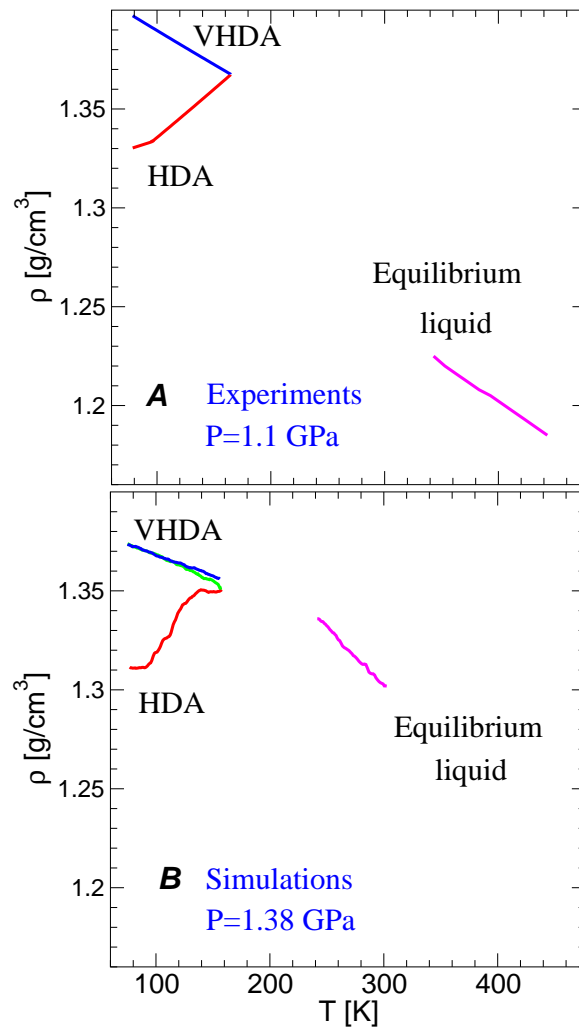


FIG. 1: Experimental data (from Ref. [9]) and simulation results for the temperature dependence of the density ρ during the conversion of HDA to VHDA at high pressure (red line) and for the cycling of VHDA between 165 and 77 K (green and blue lines). The equilibrium ρ data in the liquid state (magenta line) from the HGK equation of state [26] and for SPC/E potential are also shown. Numerical results for different pressure values are qualitatively similar to the results shown in the figure for $P=1.38$ GPa. The density difference between HDA and VHDA decreases on increasing P . We also find that the density dependence of the VHDA recovered at $T=77$ K and ambient pressure varies between 1.22 and 1.28 g/cm³.

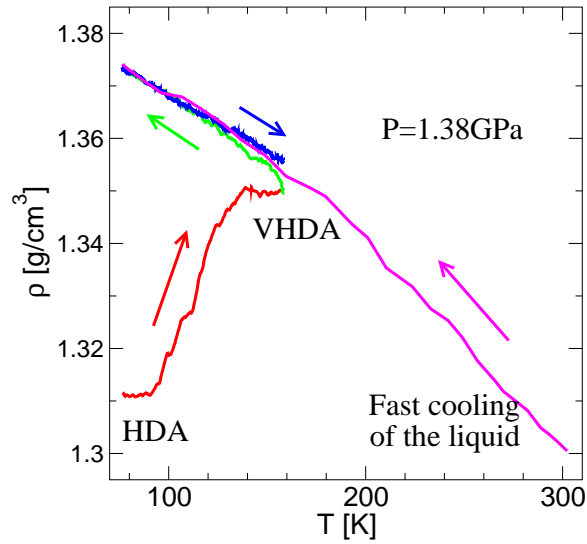


FIG. 2: Comparison between the temperature-dependence of the density in VHDA (from Fig.1B) and the temperature-dependence of the density during continuous cooling of the liquid down to 77 K, at a quenching rate of -10^4 K/ns.

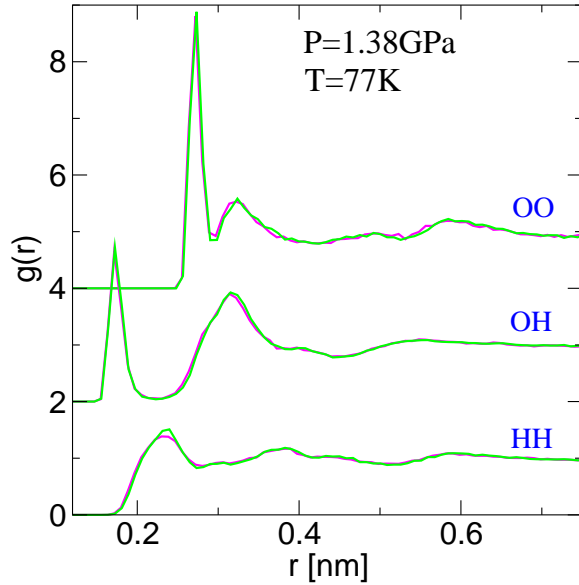


FIG. 3: Comparison between the structural properties of VHDA (green) and the glass obtained by cooling the liquid under pressure calculated at $T = 77$ K and $P = 1.38$ GPa (magenta). Both oxygen-oxygen, oxygen-hydrogen, hydrogen-hydrogen radial distribution functions are shown. Despite the extremely different previous histories, the VHDA structure appears identical to the structure of the liquid cooled under pressure.

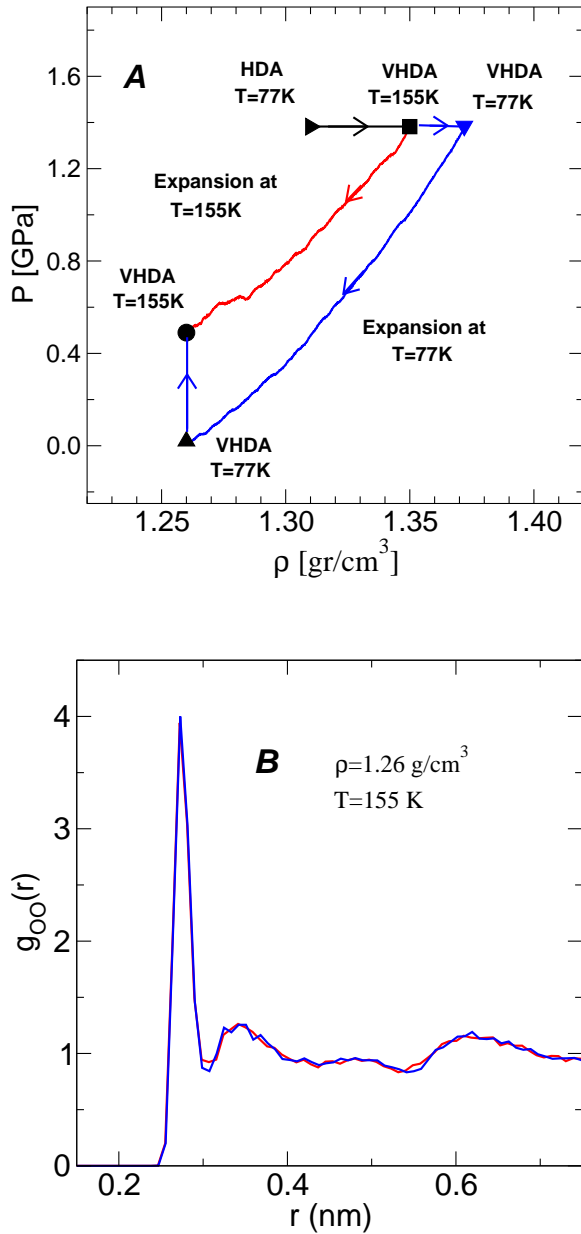


FIG. 4: Demonstration that at $T = 155$ K and $\rho = 1.26$ g/cm³ two glasses, each with completely different histories (red and blue curves), are identical both (A) thermodynamically and (B) structurally. (A) P – ρ diagram. Blue curve: recovering of the VHDA at ambient pressure at $T = 77$ K followed by an isochoric heating at density 1.26 g/cm³ up to $T = 155$ K. Red curve: isothermal decompression of VHDA at $T = 155$ K down to density 1.26 g/cm³. The final product is independent of the path despite the different histories. The important point is that no transformation from VHDA to HDA is observed. (B) Comparison between the radial distribution function of the two glasses obtained following the blue and the red paths in (A).